

Electrostatic Barrier Against Dust Growth in Protoplanetary Disks

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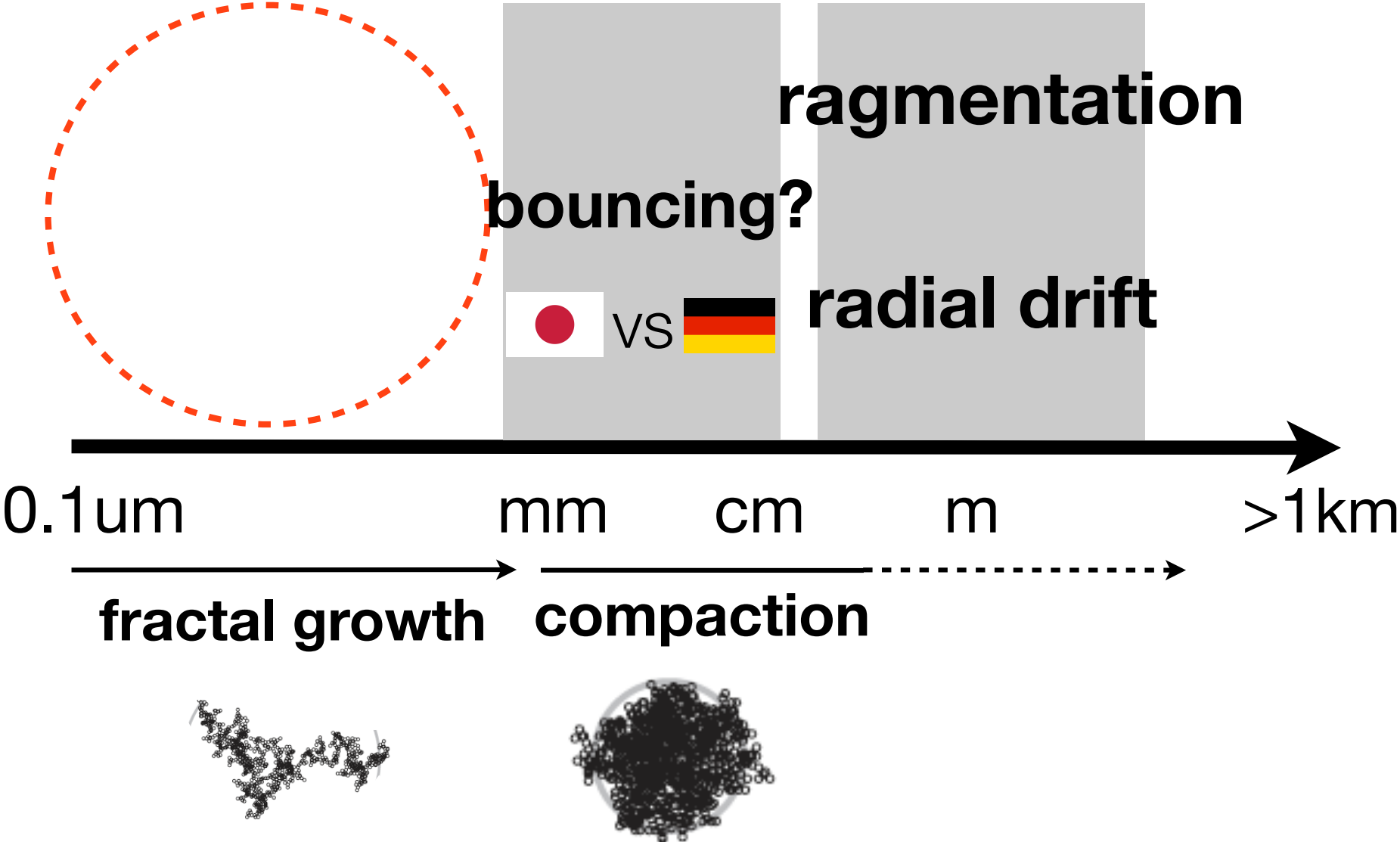
Masa-aki Sakagami (Kyoto U.)

Ref: Okuzumi 2009, ApJ, 698, 1122

Okuzumi et al. 2009, ApJ, 707, 1247

Okuzumi et al., submitted, [arXiv:1009.3199](https://arxiv.org/abs/1009.3199); [1009.3101](https://arxiv.org/abs/1009.3101),₁

Barriers Against Dust Growth in PP disks



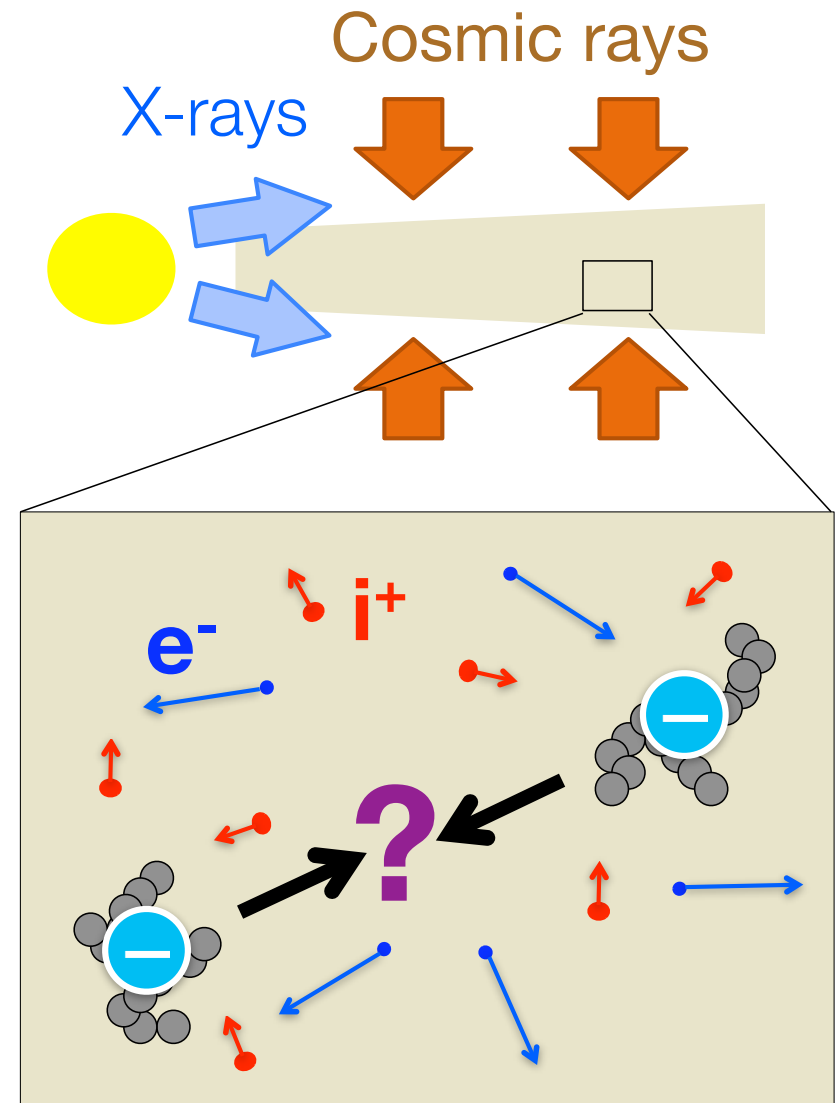
The “Charge Barrier” Against Dust Coagulation?

* PP-disks are **weakly ionized** by various high-energy sources:

- ▶ cosmic rays (e.g., Umebayashi & Nakano 80)
- ▶ stellar X-rays (e.g., Glassgold et al. 97)

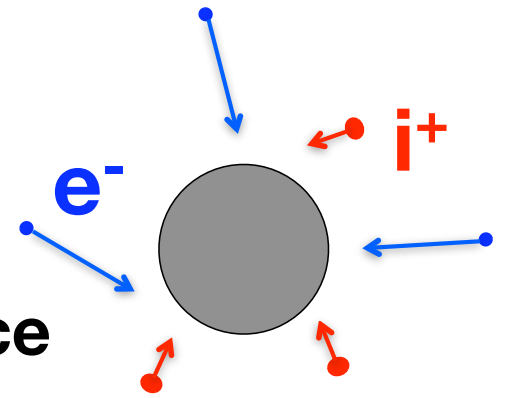
* Dust particles in an ionized gas: (on average) get **negative** charges!
(first discovered in astrophysics (Spitzer 41), now better known in plasma physics)

* “Assymmetric” ($\langle Q \rangle \neq 0$) charging
➔ Coulomb barrier between colliding particles



Prelude: Grain Charging in a *Fully* Ionized Gas

- number density: $n_i \approx n_e$ ($\ll |Q|n_d$)
- thermal speed: $v_i \ll v_e$ ($\leftarrow m_i \gg m_e$)
- ➔ incident flux: $n_i v_i \ll n_e v_e$
- ➔ the grain charges up until the Coulomb force equilibrates these fluxes.



Equilibrium Condition:

$$-eV \sim k_B T$$

$V = Q/a$: surface potential of the grain (e.g., Spitzer 1941)

➔ $V \sim -\frac{k_B T}{e} \sim -10 \text{ mV} \left(\frac{T}{10^2 \text{ K}} \right)$ ← independent of grain size a

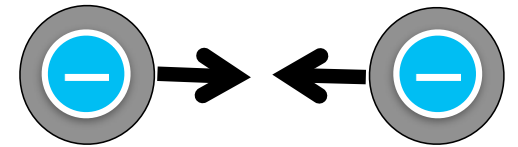
➔ $Q = aV \sim -\text{a few } e \left(\frac{a}{0.1 \mu\text{m}} \right) \left(\frac{T}{10^2 \text{ K}} \right) \propto a$

Prelude: Collision of Charged Grains

Consider a collision between two charged grains

Hitting Condition: $E_{\text{kin}} > E_{\text{el}}$

- $E_{\text{kin}} = \frac{1}{2} M_{\text{red}} (\Delta v)^2$: kinetic energy
- $E_{\text{el}} = \frac{Q_1 Q_2}{a_1 + a_2}$: electrostatic energy



Using $Q = aV$, $E_{\text{el}} \sim aV^2 \sim \left(\frac{a}{0.1 \mu\text{m}}\right) \left(\frac{T}{10^2 \text{ K}}\right) k_{\text{B}} T \propto a$

If the relative velocity is driven **only** by Brownian motion ($E_{\text{kin}} \sim k_{\text{B}} T$),
collision is severely inhibited at **$a \gg 0.1 \mu\text{m}!!$**

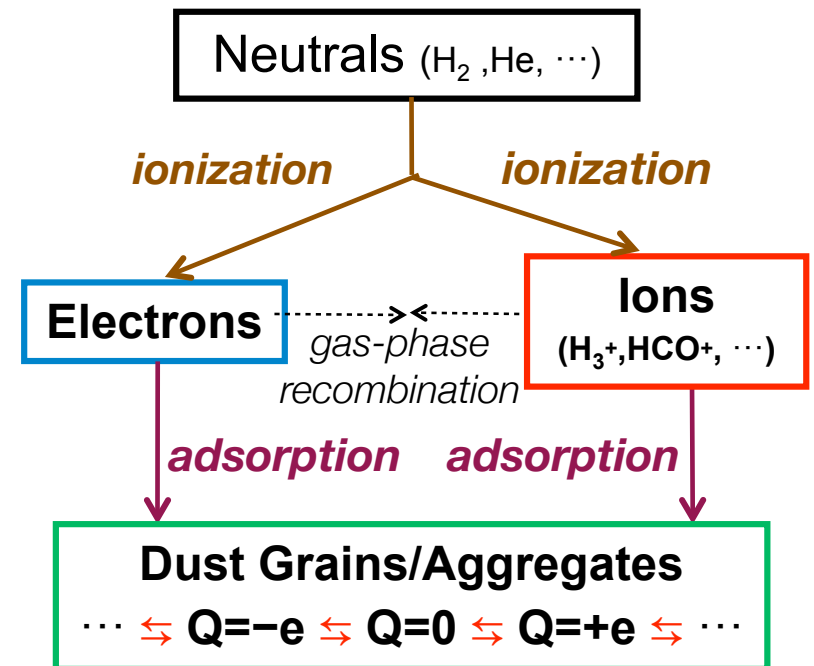
This Work

- Dust Charging in **weakly ionized** gases
- The role of the size distribution
- The effect of settling/turbulence

Dust Charging in Weakly Ionized Gases

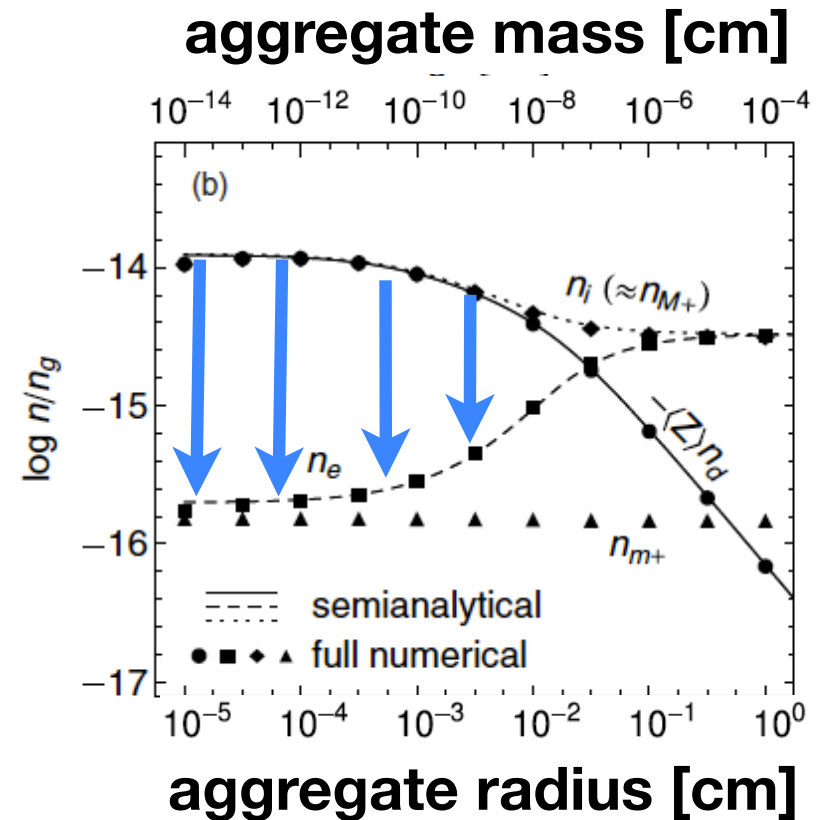
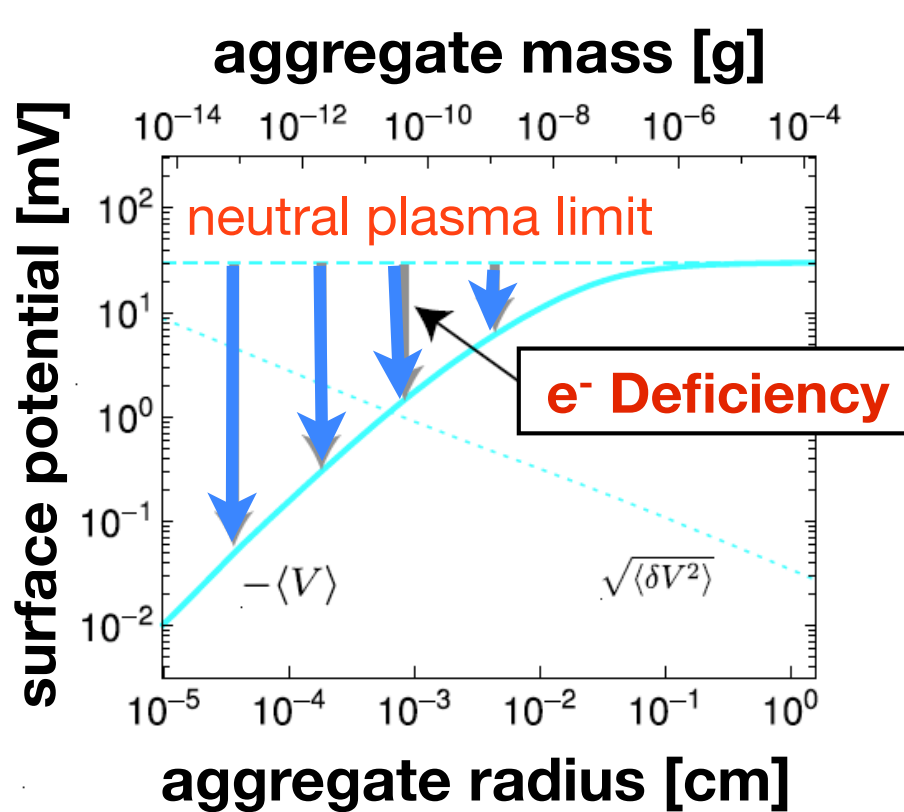
- * Dust charge distribution $n_d(Z)$
- * Ion & electron number density n_i, n_e

are mutually dependent!
→ **Must be computed simultaneously!**



Electron Deficiency Effect (Okuzumi 2009)

MMSN + cosmic rays + X-rays, $r=5\text{AU}$, $z=H$
 Fractal growth, $D=2$ (0.1 μm monomer)



Weak ionization prevents small aggregates from being strongly charged

Coagulation of Charged Dust: Simulations

▶ **Smoluchowski's method** to follow the evolution of size distribution

▶ **Collisional cross section** *including Coulomb correction*:

$$\sigma_{\text{eff}} = \pi(a_1 + a_2)^2 \left(1 - \frac{E_{\text{el}}}{E_{\text{kin}}} \right) \quad E_{\text{kin}} = \frac{1}{2} \frac{M_1 M_2}{M_1 + M_2} \Delta v^2 \quad E_{\text{el}} = \frac{Q_1 Q_2}{a_1 + a_2} = \frac{a_1 a_2}{a_1 + a_2} V^2$$

* Δv : **Relative velocity** = Brownian motion + **settling** + **turbulence**

* V : **Dust surface potential** ← analytic solution by Okuzumi (2009)

▶ **“Hit-and-stick” (fractal) aggregation model** (Okuzumi et al. 2009)

- no compaction nor fragmentation

- determines the porosity of collision products using an empirical formula obtained from N-body calculations

▶ **Local (0-dim.) simulation** (advection is neglected)

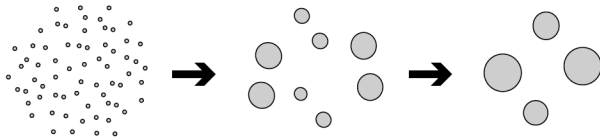
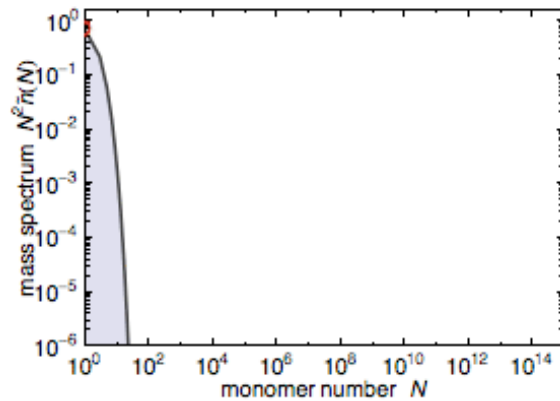
Key parameters: ionization rate ζ , “drift acceleration” g

The Outcome of Dust Growth

Three types of outcomes depending on ζ & g :

(Okuzumi et al. submitted)

“Orderly” growth



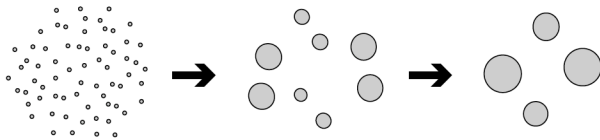
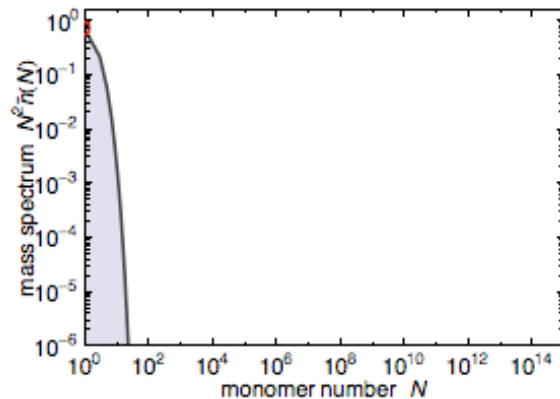
- ▶ No-charging cases
- ▶ Relatively narrow size distribution

The Outcome of Dust Growth

Three types of outcomes depending on ζ & g :

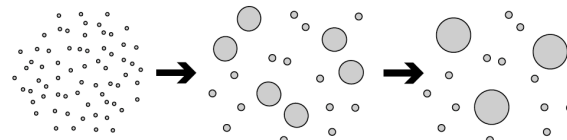
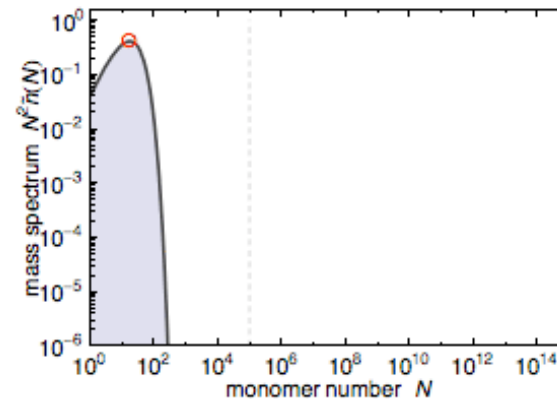
(Okuzumi et al. 2010a, submitted)

“Orderly” growth



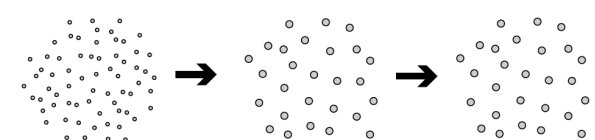
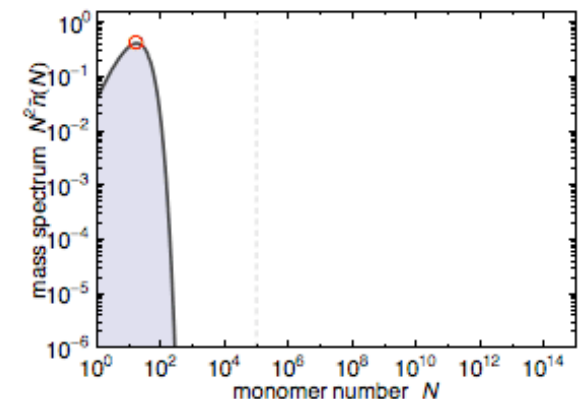
- ▶ When g is high
- ▶ Charge does not affect the growth

“Bimodal” growth



- ▶ Low g but low ζ
- ▶ *Electron deficiency effect!*

“Freezeout”



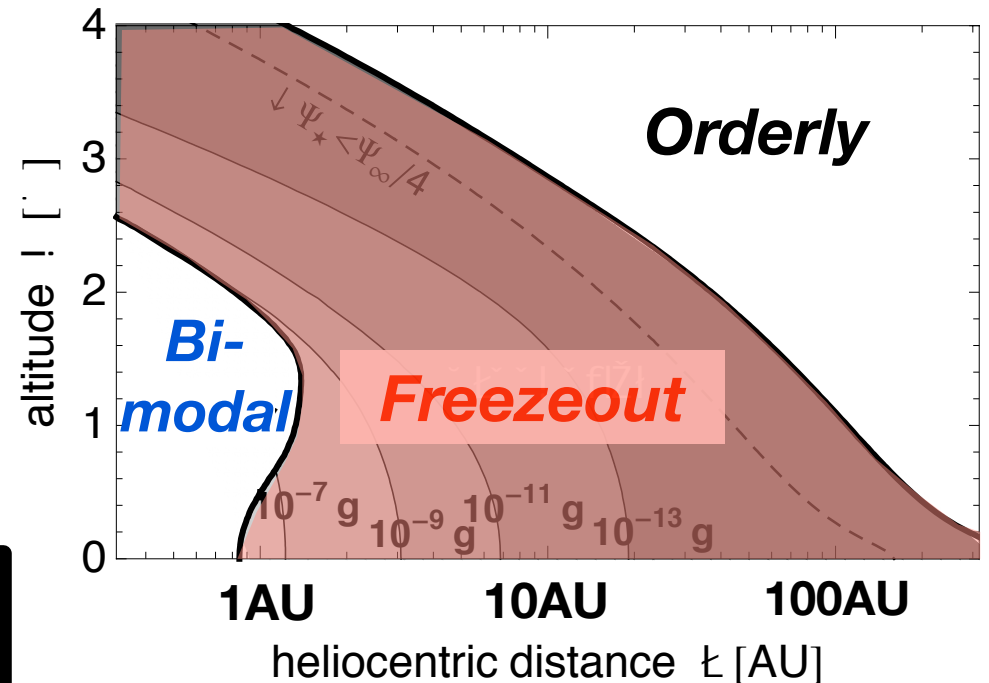
- ▶ Low g & High ζ
- ▶ Growth stops in the Brownian motion regime

Application to Protoplanetary Dust Growth

- ✓ MMSN model (Hayashi 81), isothermal
- ✓ Cosmic rays, stellar X-rays, ^{26}Al
- ✓ Monomer size: 0.1 μm
- ✓ Includes turbulence-driven collision velocity
(but *ignores* turbulent diffusion)

- ▶ **At depth $\leq 0.1 \text{ g/cm}^2$, orderly growth**
 - ← High collision velocity ← Low gas density
- ▶ **At depth $\geq 10^3 \text{ g/cm}^2$, bimodal growth**
 - ← Electron deficiency ← Weak ionization
- ▶ **At intermediate depths, freezeout!**

(Okuzumi et al. 2010b, submitted)

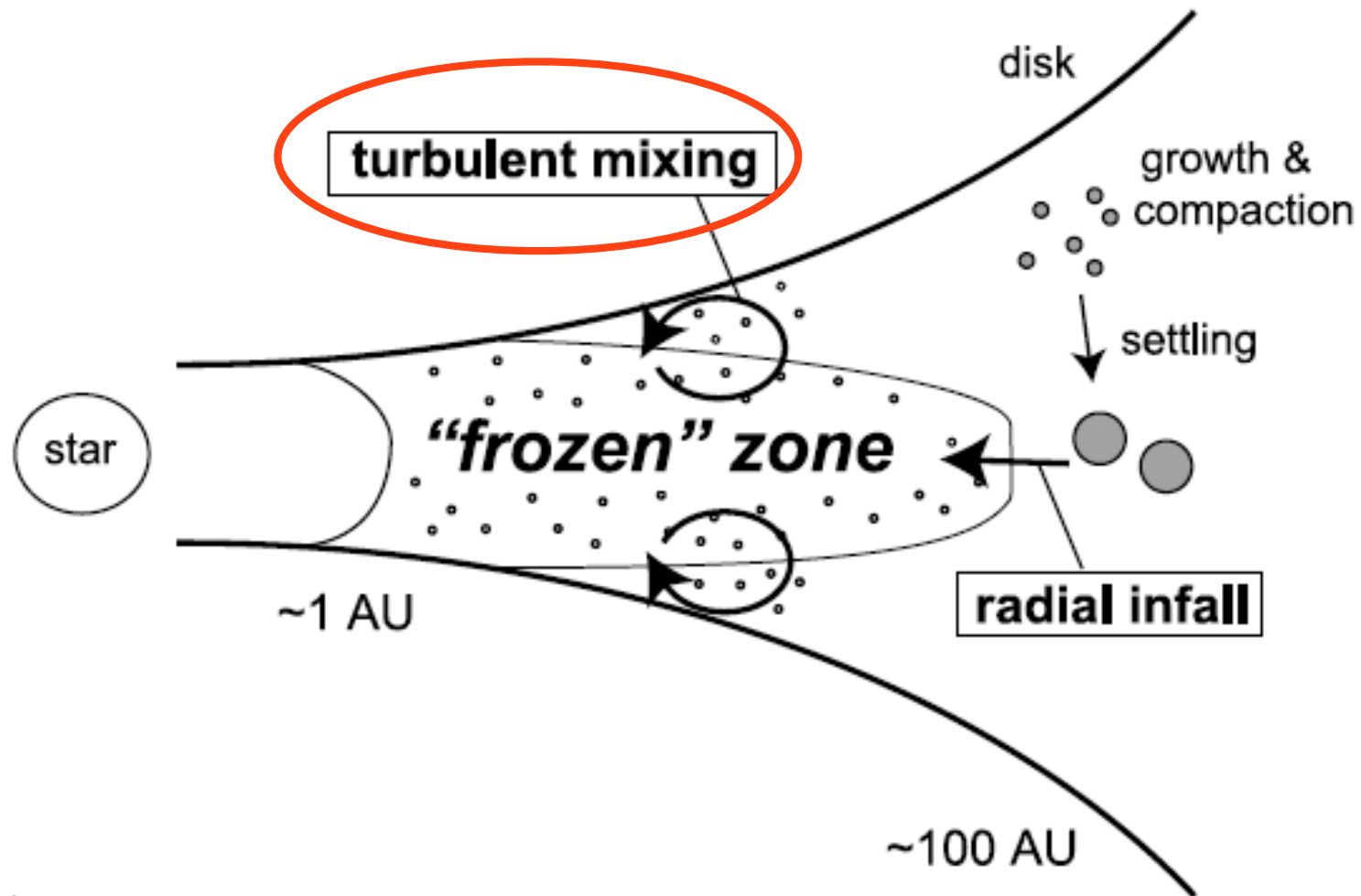


The maximum size of the
“frozen” aggregates in the disk

10^{-7} g (0.1mm, 10^{-3} g/cc) !!

$(\ll M_{\text{comp}} \sim 10^{-4} \text{ g})$

Dust Transport Across Frozen-zone Boundaries



(Okuzumi et al. 2010b, submitted)

Dust Growth Timescales

Mean Collision Time at distance r

$$\tau_{\text{grow}} \equiv \left(\frac{d \ln M}{dt} \right)^{-1}$$

W/o charging, $\tau_{\text{grow}} \sim \frac{\Sigma_g}{\Sigma_d} T_K$

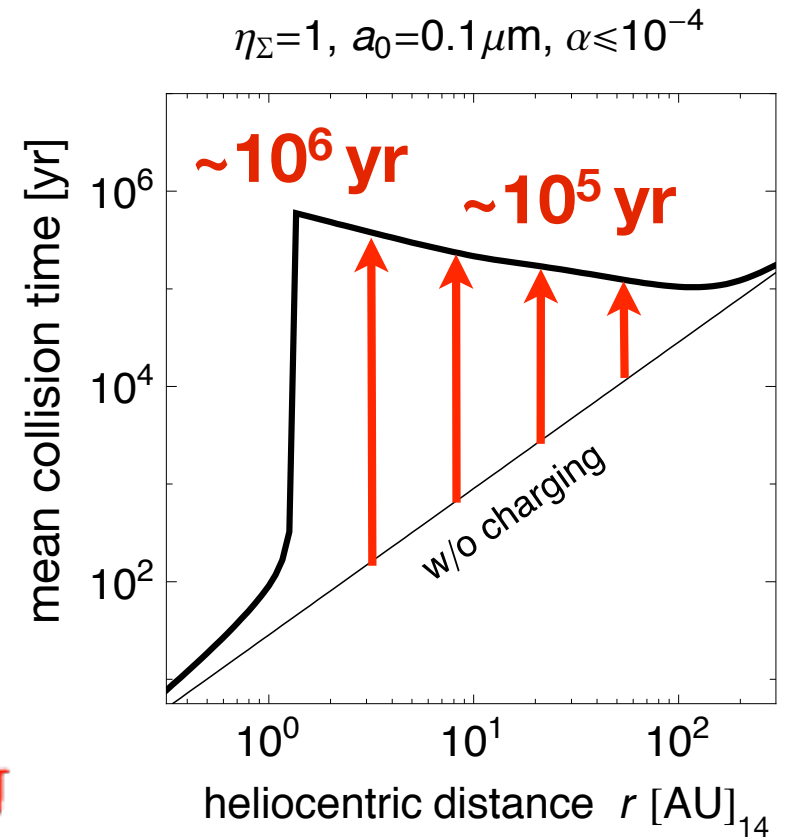
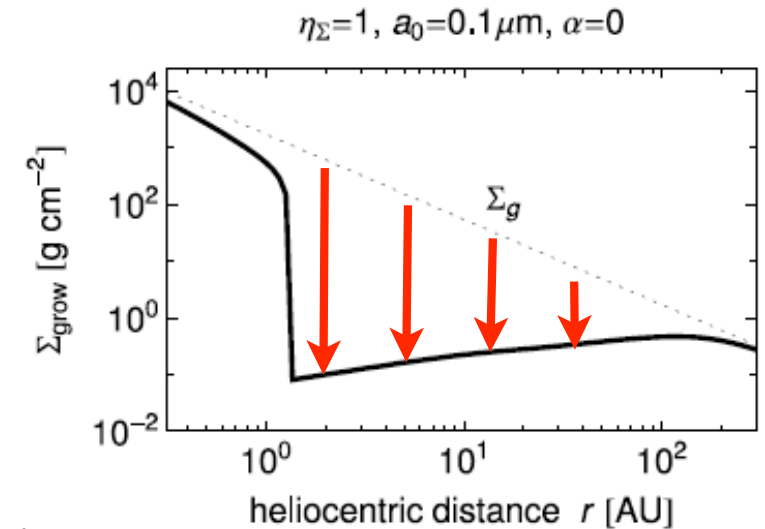
$\sim 10^2 \text{ yr @1AU}$

$\sim 10^5 \text{ yr @100AU}$

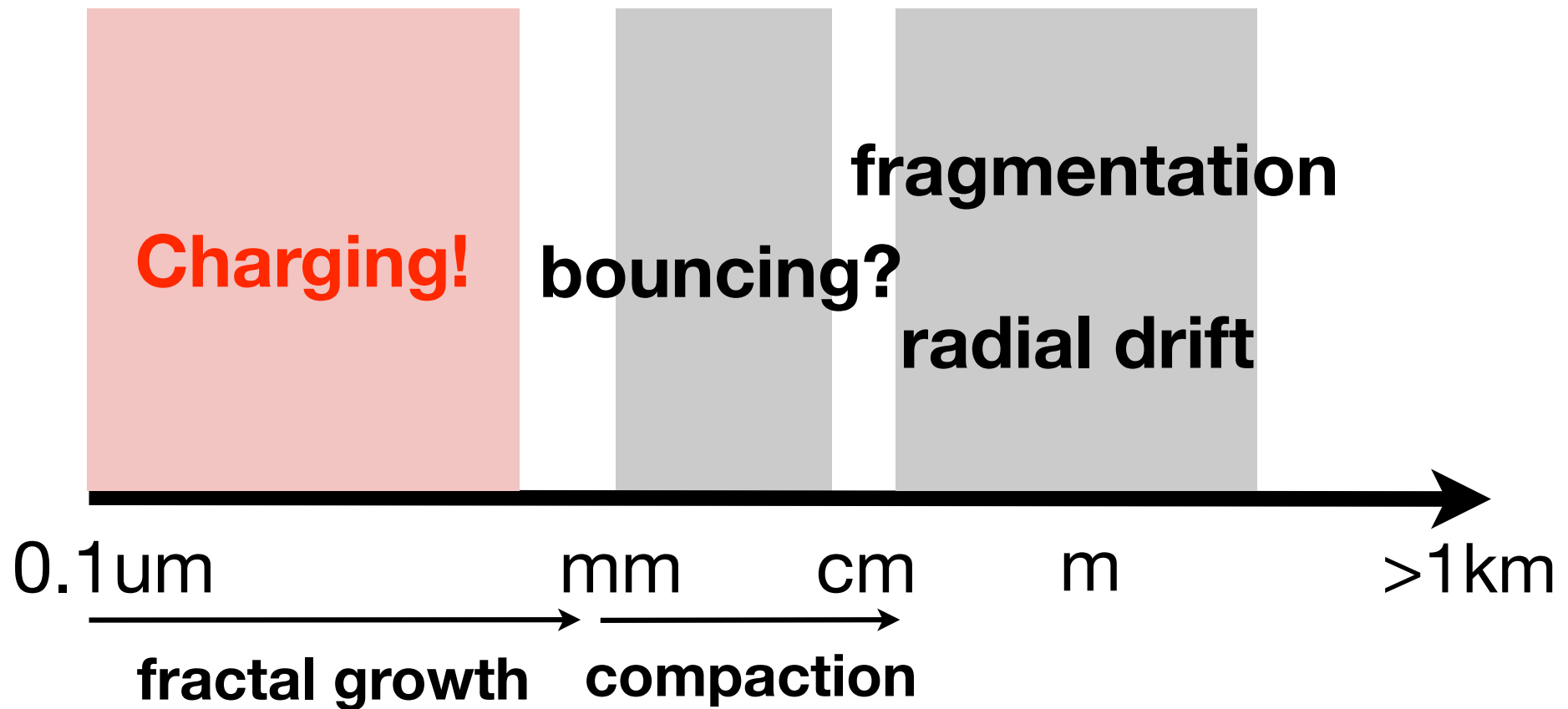
W/ charging, $\tau_{\text{grow}} \sim \frac{\Sigma_g}{\Sigma_{\text{grow}}} \frac{\Sigma_g}{\Sigma_d} T_K$

$\sim 10^6 \text{ yr @1AU}$

$\sim 10^5 \text{ yr @100AU}$



Conclusion



- Charging prevents dust growth **before compaction occurs!**
- At least, the “frozen” tiny aggregates are retained in the disk **over a timescale of 10^6 years!**