

Staub in Planetensystemen/惑星系の「うちゅうじん」

Sep. 27 - Oct. 1, 2010, Jena, Germany

Formation of cosmic crystals by eccentric planetesimals

***H. Miura¹, K. K. Tanaka², T. Yamamoto², T. Nakamoto³,
J. Yamada¹, K. Tsukamoto¹, and J. Nozawa¹***

¹Tohoku Univ., Japan

²Hokkaido Univ., Japan

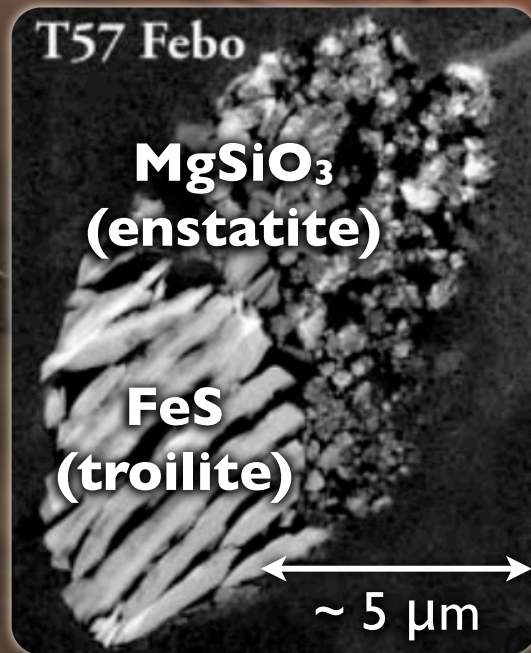
³Tokyo Inst. of Tech., Japan

**This study has been already published.
(Miura+2010, ApJ 719, 642-654)**



COSMIC CRYSTALS

Fine dust from comet



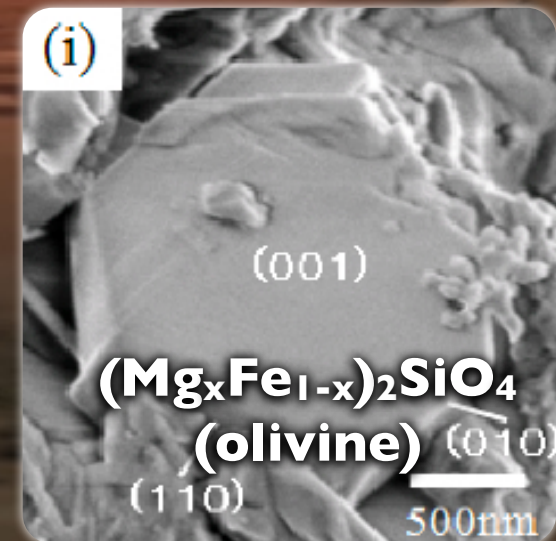
Stardust mission
(e.g., Brownlee+2006,
Science 314, 1711)

IDPs



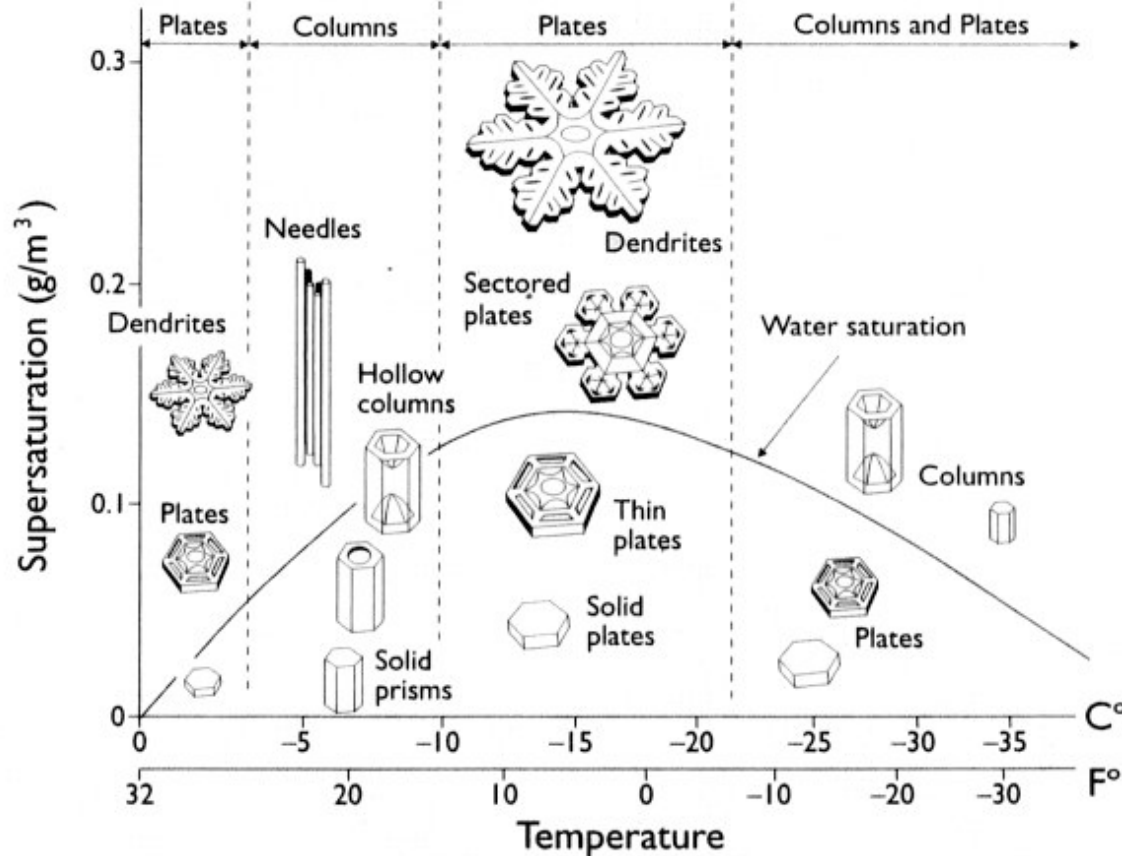
(Bradley+1983,
Nature 301, 473)

Primitive meteorite

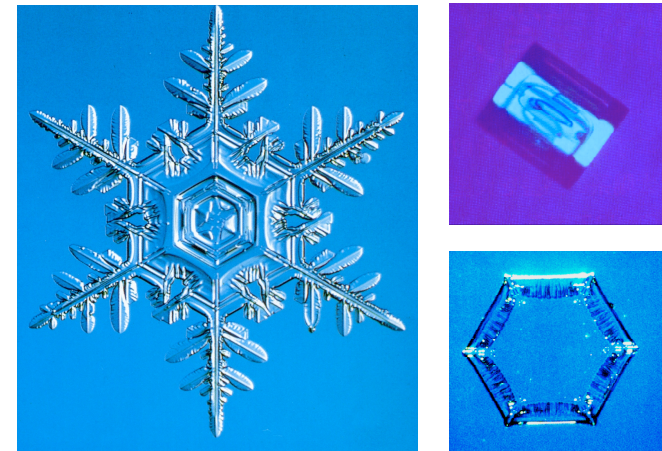


(Nozawa+2009,
Icarus 204, 681)

Introduction: Morphodrom (snowflake)



Snowflake changes its shape depending on temperature (undercooling) and supersaturation (density of water vapor)
(Nakaya diagram)



From Prof. Furukawa, Hokkaido Univ.
<http://www.lowtem.hokudai.ac.jp/ptdice/>

Morphologies of crystals reflect their formation condition

relationship
morphologies v.s.
formation condition

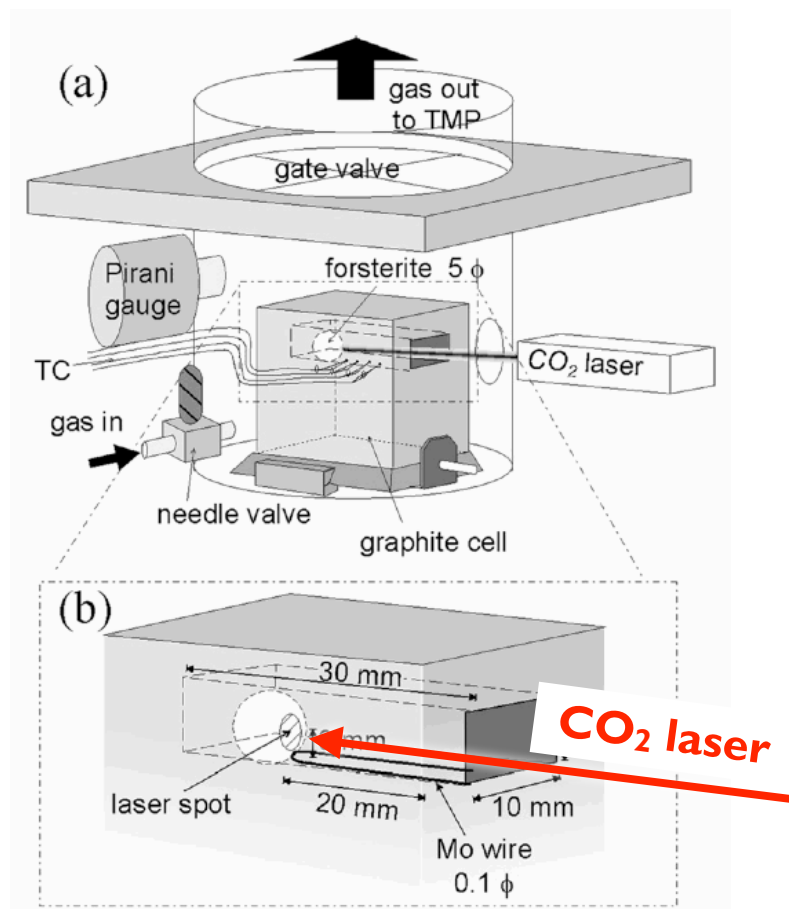


“morphodrom”

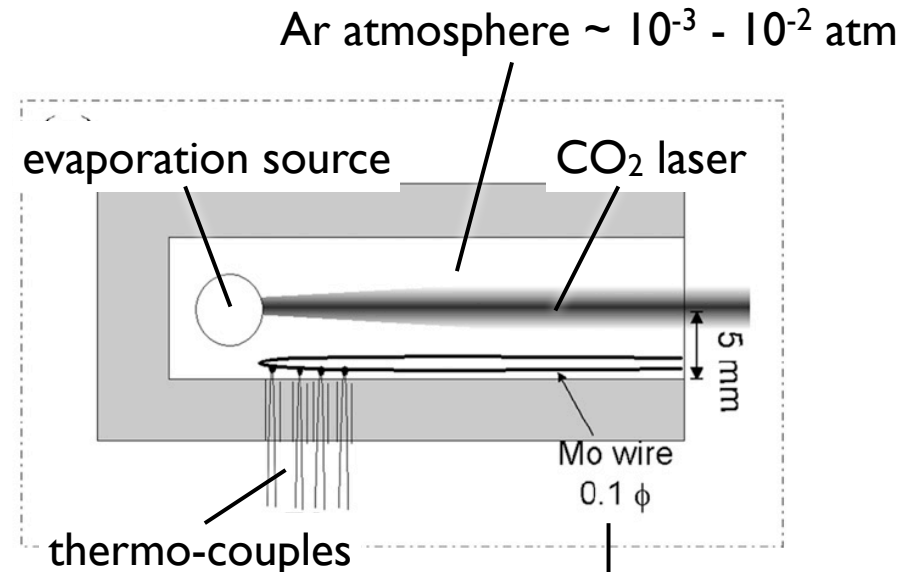
Introduction:

Evaporation + condensation experiments

flash heating by
CO₂ laser irradiation
(Kobatake+2008, Icarus 198, 208)



top view of graphite cell

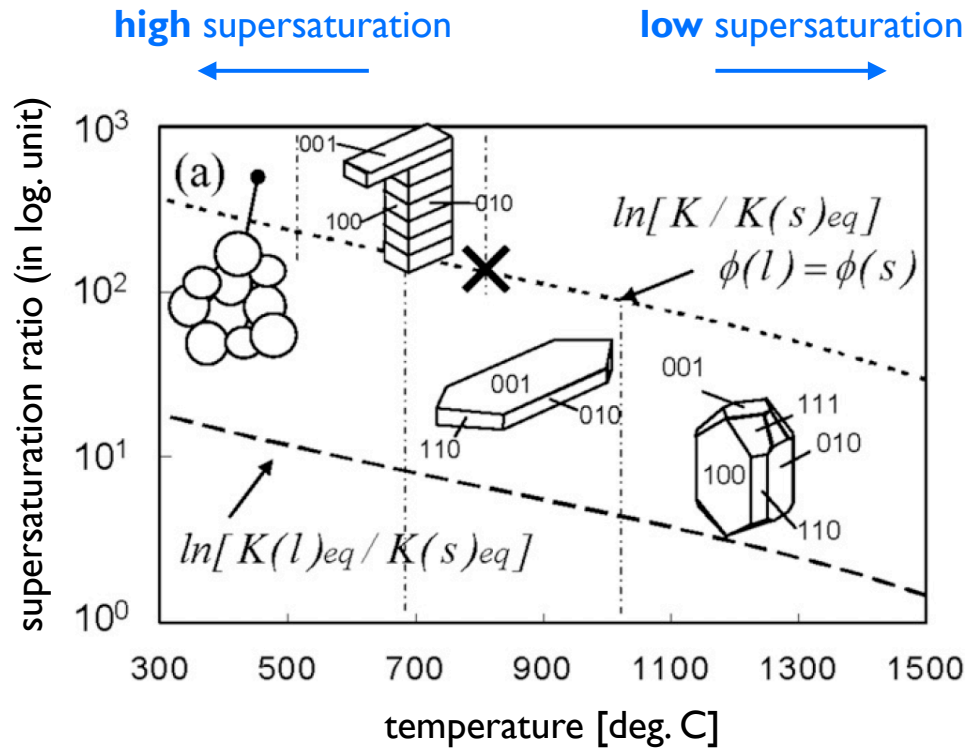


temperature
measurement

sample correction

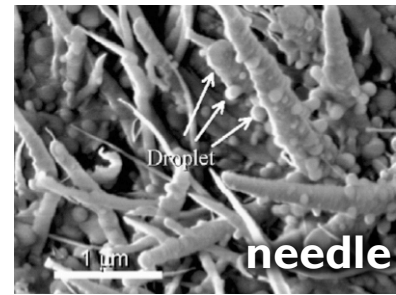
- chemical composition (EDS)
- morphology (FE-SEM)
- identification (TEM)

Introduction: Morphodrom (forsterite)

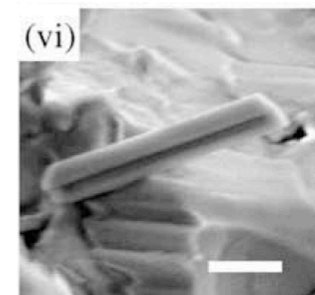


similarity in morphology
↓
vapor growth
at high supersaturation

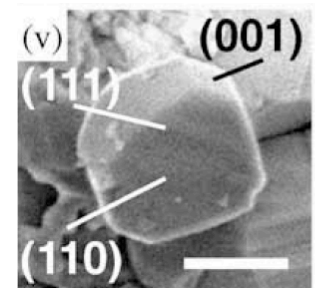
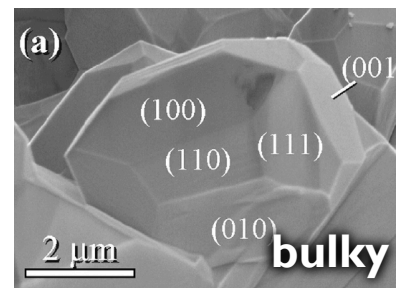
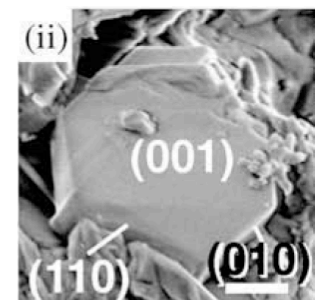
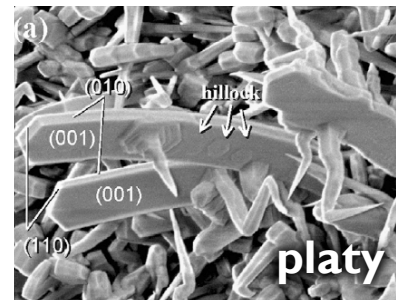
synthesized
(Kobatake+2008, Icarus 198, 208)



in matrix of Allende
(Nozawa+2009, Icarus 204, 681)

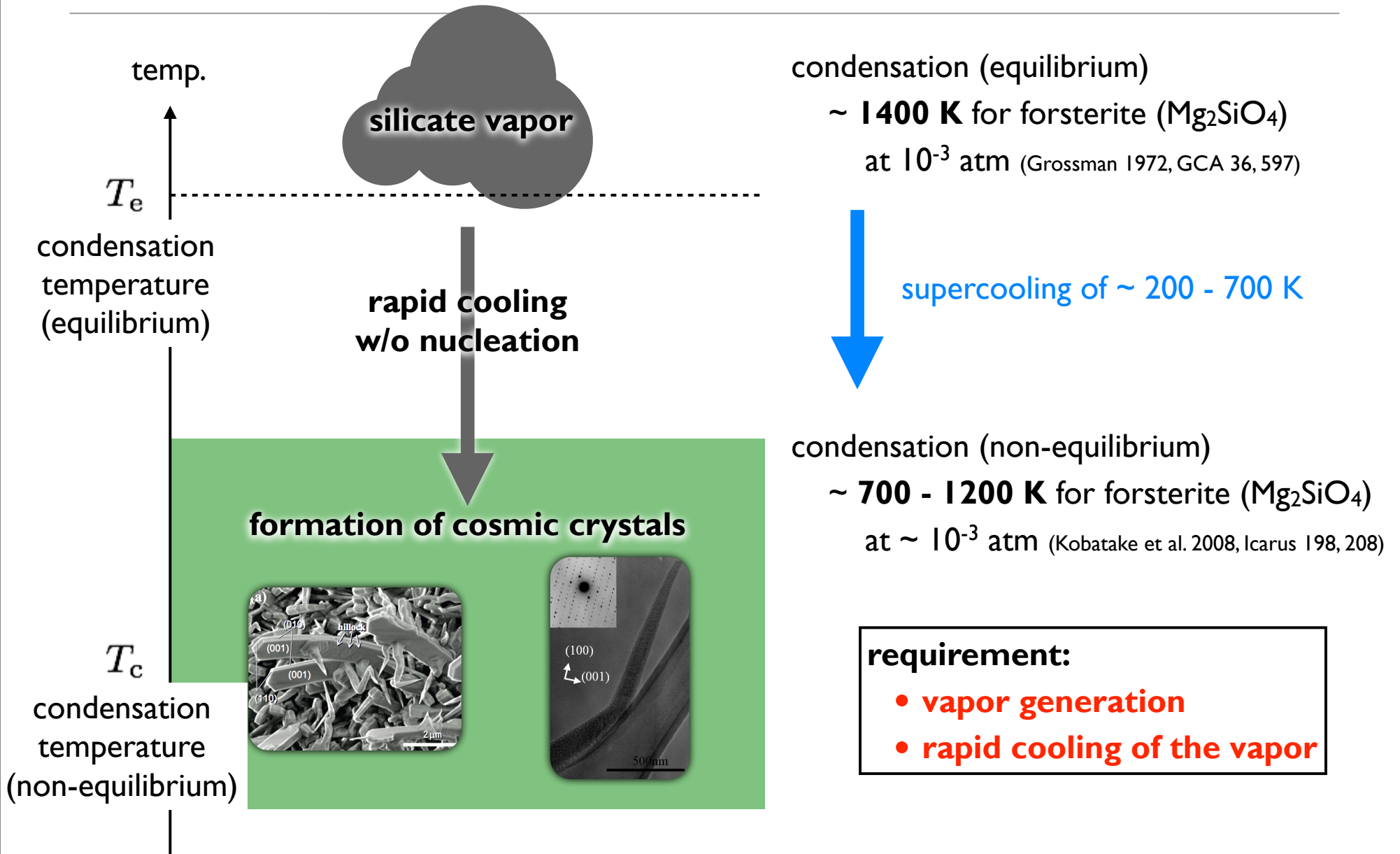


↑
supersaturation



Introduction:

Condensation in non-equilibrium



Introduction:

Candidate of cosmic crystal formation

eccentric planetesimals:

During planet formation, planetesimals take eccentric orbits because of gravitational interaction between themselves.

- **Jovian resonances**

(Weidenschilling+1998, Science 279, 681)

- **Dynamical shake-up**

(Nagasawa+2005, ApJ 635, 578)

relative velocity (supersonic) between

- nebular gas ($e \sim 0$)
- planetesimals ($e > 0$)



“bow shock”

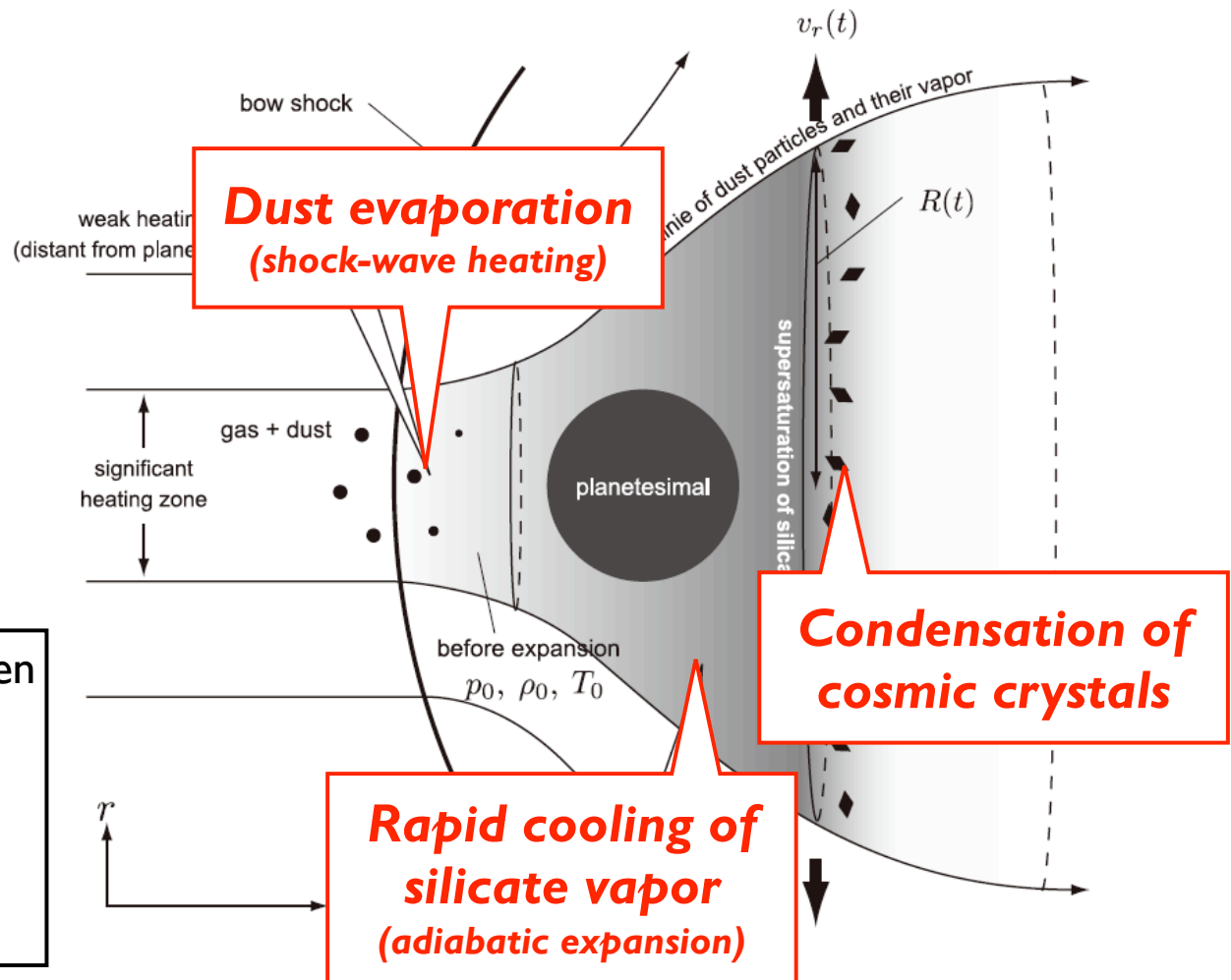


Fig. schematic of planetesimal bow shock

(Miura+2010, ApJ 719, 642)

I. Dust evaporation

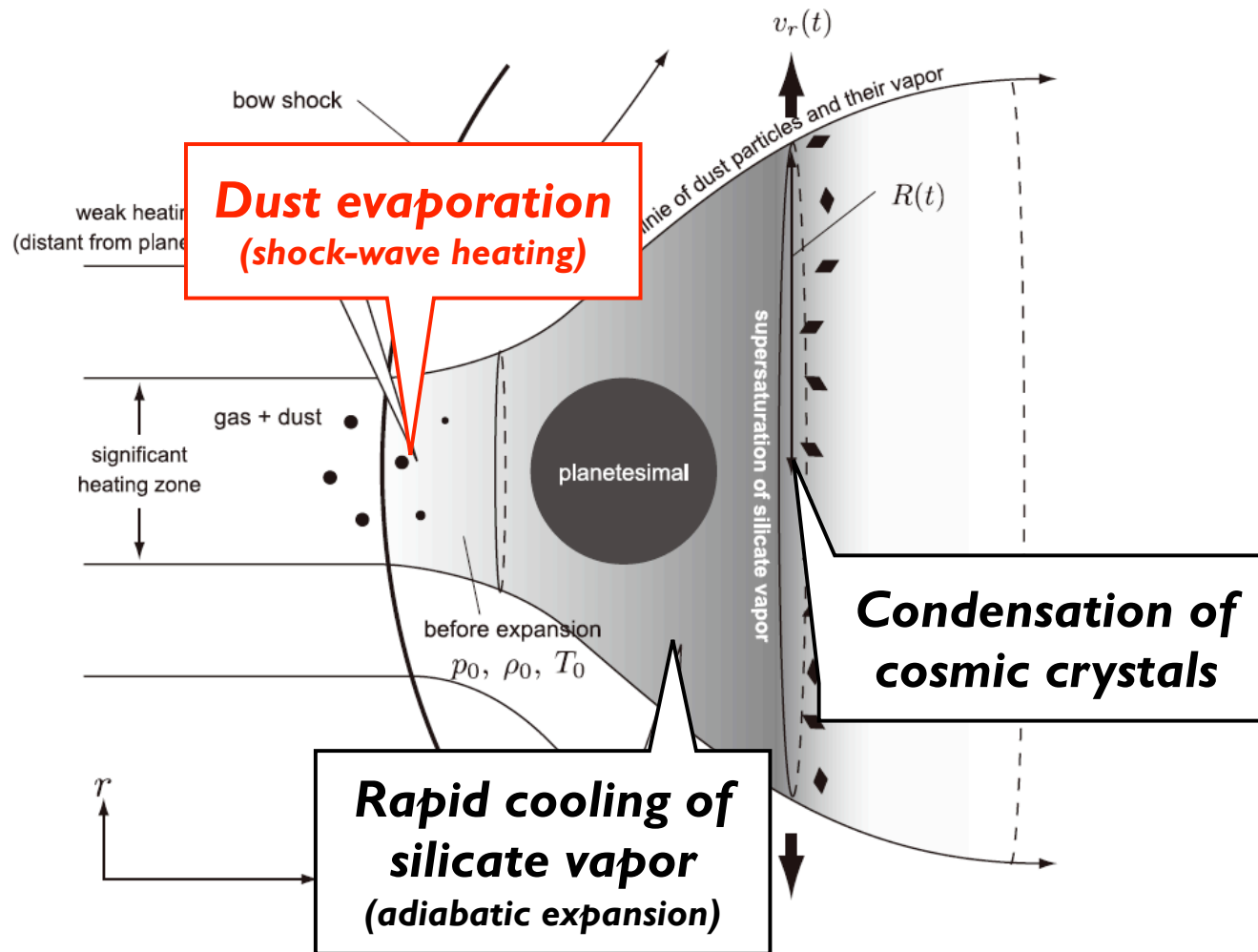
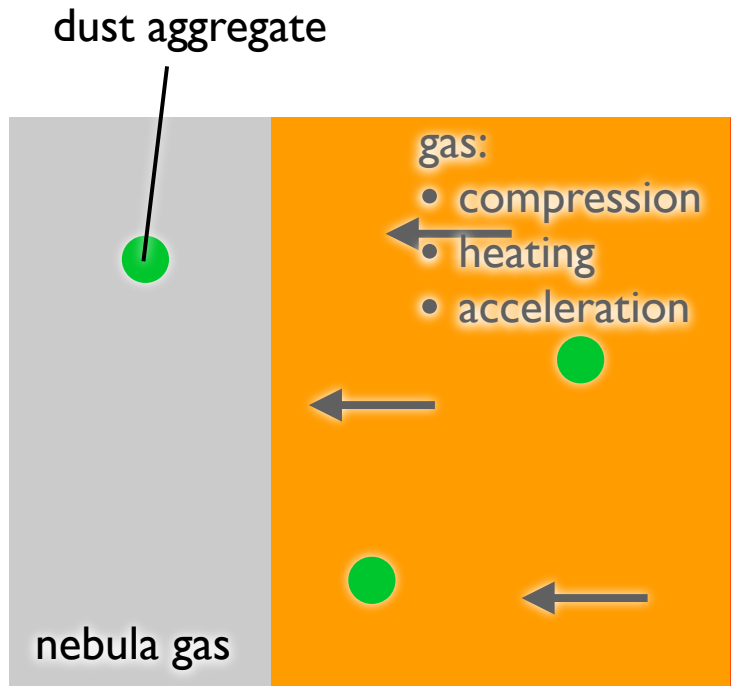


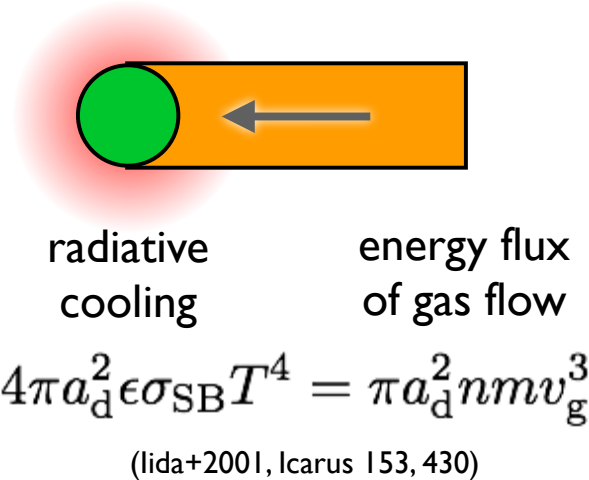
Fig. schematic of planetesimal bow shock
(Miura+2010,ApJ 719, 642)

I. Dust evaporation

Shock-wave heating

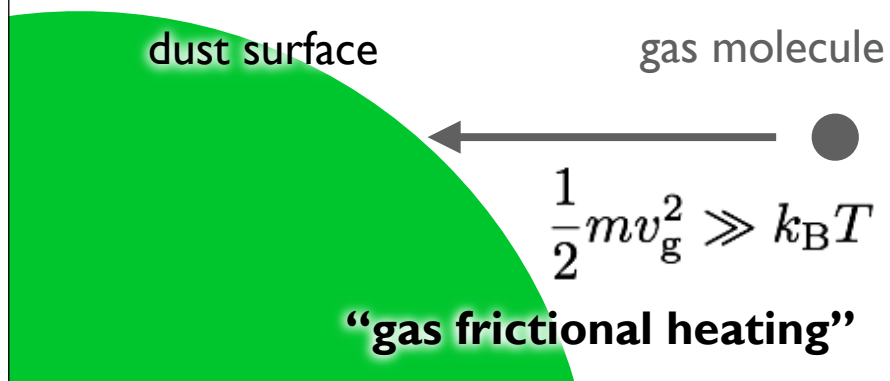


How strong is the gas frictional heating?



For $v_g = 10 \text{ km s}^{-1}$ and $n = 10^{15} \text{ cm}^{-3}$, we obtain

$$T_{\text{peak}} \sim 1720 \text{ K}$$

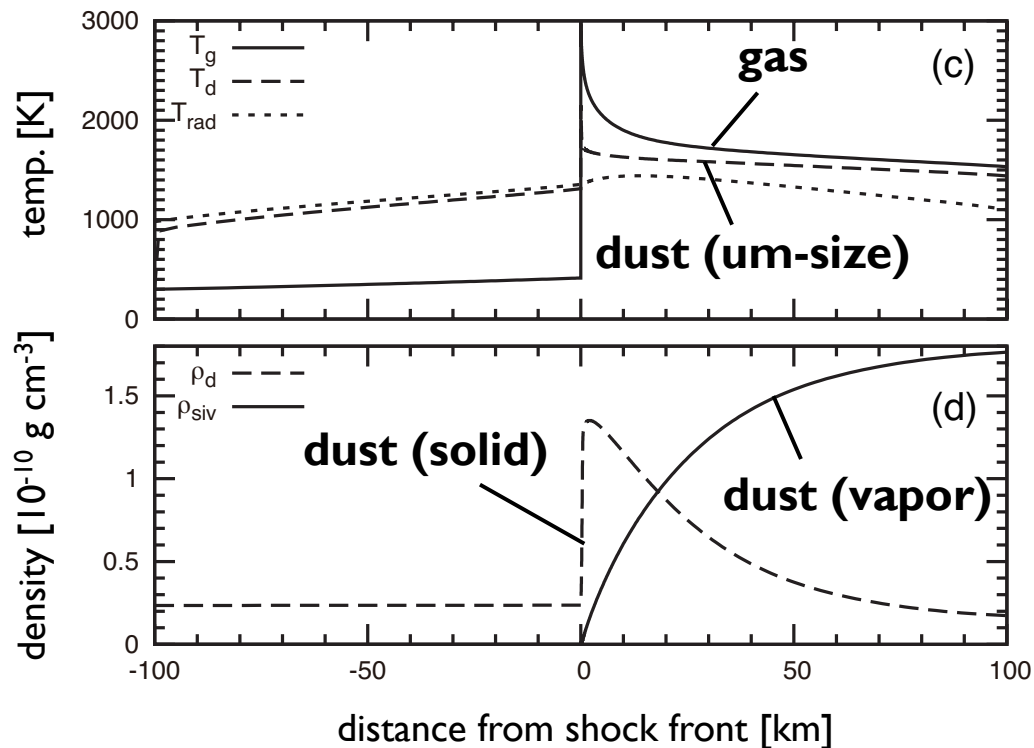


- melting of silicate dust aggregates (chondrule formation)
(Wood 1984, EPSL 70, 11 and others)
- evaporation of μm -sized particles
(Miura+2005, Icarus 175, 289)

I. Dust evaporation

Dust in hot gas

1D plane-parallel steady model
(Miura+2010, ApJ 719, 642)



input parameters:

planetesimal radius	$R_p = 100 \text{ km}$
gas number density (pre-shock)	$n_0 = 10^{15} \text{ cm}^{-3}$
shock velocity	$v_s = 8 \text{ km s}^{-1}$
gas/dust mass ratio	$\xi = 0.01$
dust radius	$a_d = 1 \mu\text{m}$

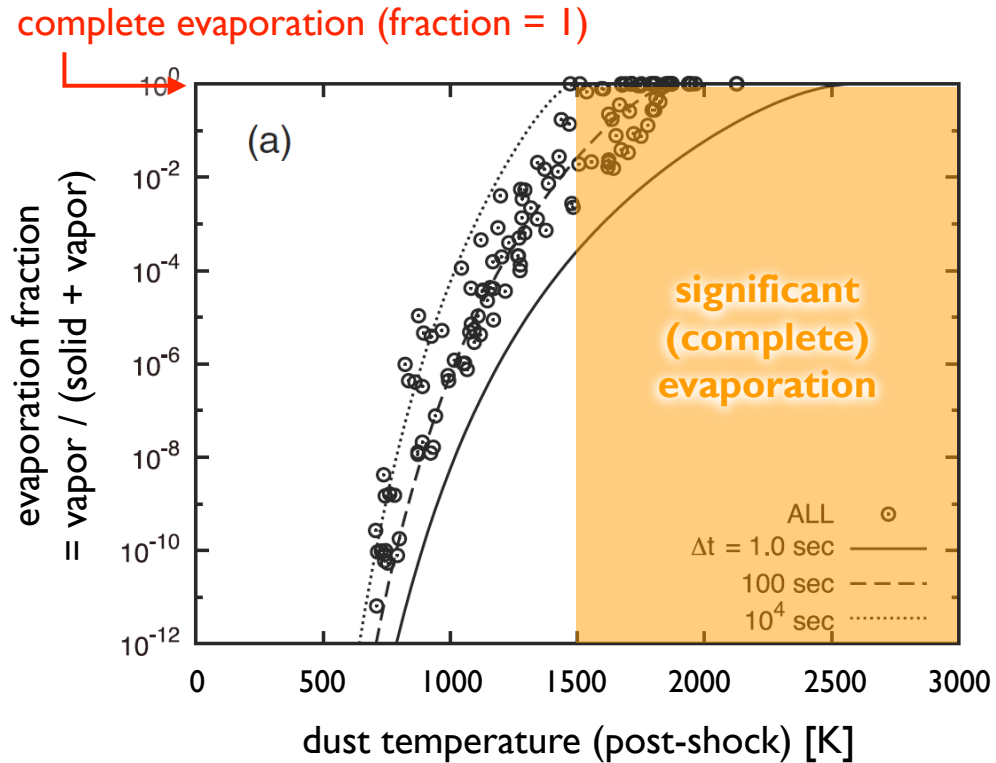
post-shock gas (far from shock front):

- temp. $\sim 1700 \text{ K}$
- density $\sim 4 \times 10^{-8} \text{ g cm}^{-3}$
- no relative velocity to dust

- dust temperature $> 1500 \text{ K}$
- evaporate significantly (90% in mass evaporates away, in this case)

I. Dust evaporation

Evaporation fraction



input parameters:

planetesimal radius	$R_p = 1 - 1000 \text{ km}$
gas number density (pre-shock)	$n_0 = 10^{13} - 10^{15} \text{ cm}^{-3}$
shock velocity	$v_s = 5 - 60 \text{ km s}^{-1}$
gas/dust mass ratio	$\eta = 0.01 - 0.1$
dust radius	$a_d = 1 \text{ }\mu\text{m}$

“significant vapor generation by planetesimal bow shock”

2. Rapid cooling of silicate vapor

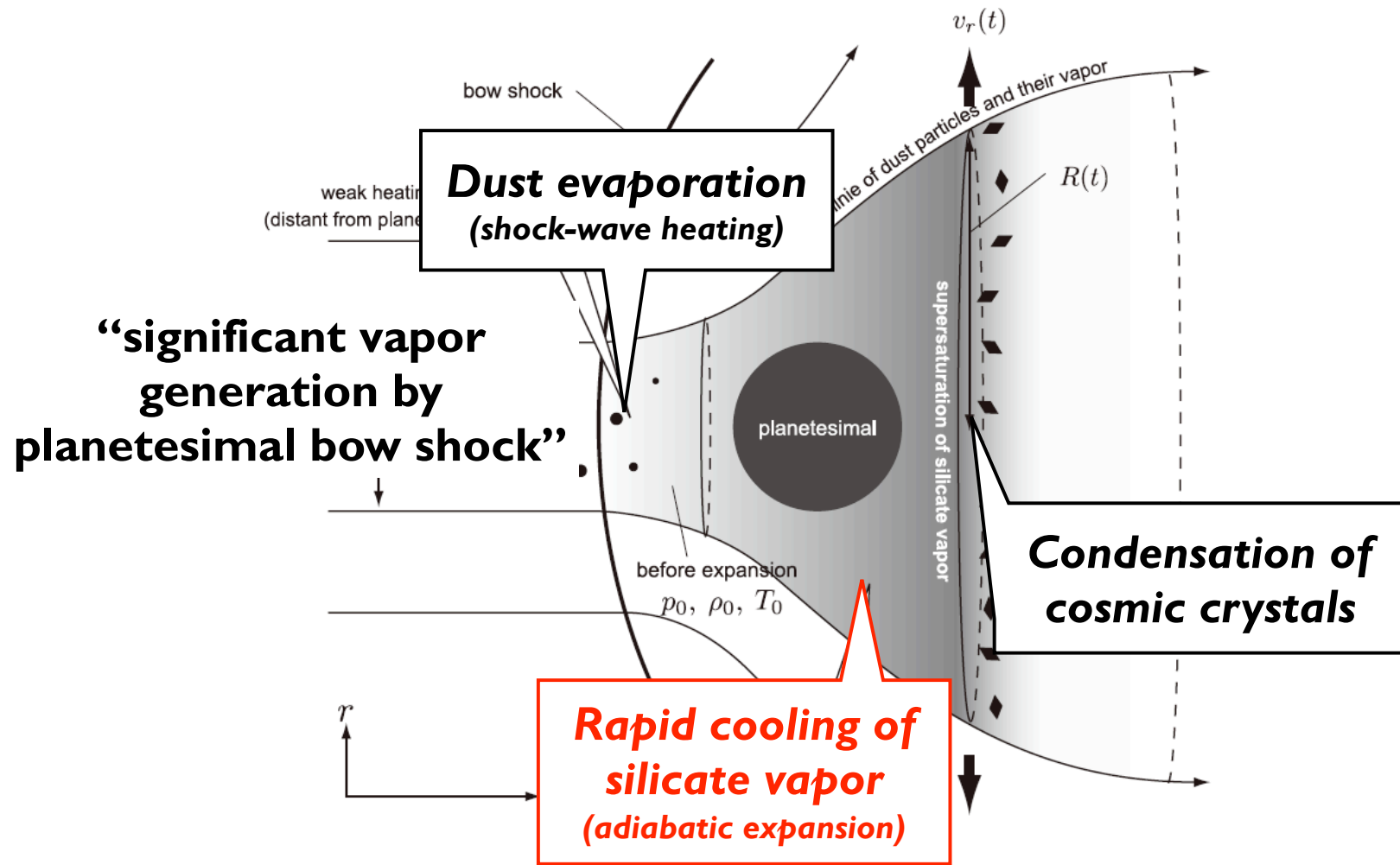


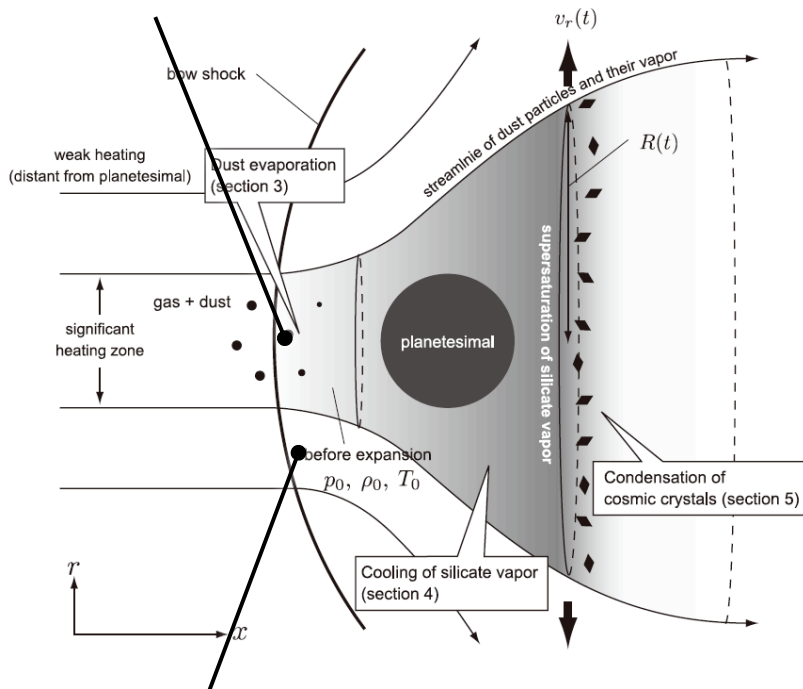
Fig. schematic of planetesimal bow shock
(Miura+2010,ApJ 719, 642)

2. Rapid cooling of silicate vapor

Expansion of shocked gas

vertical shock

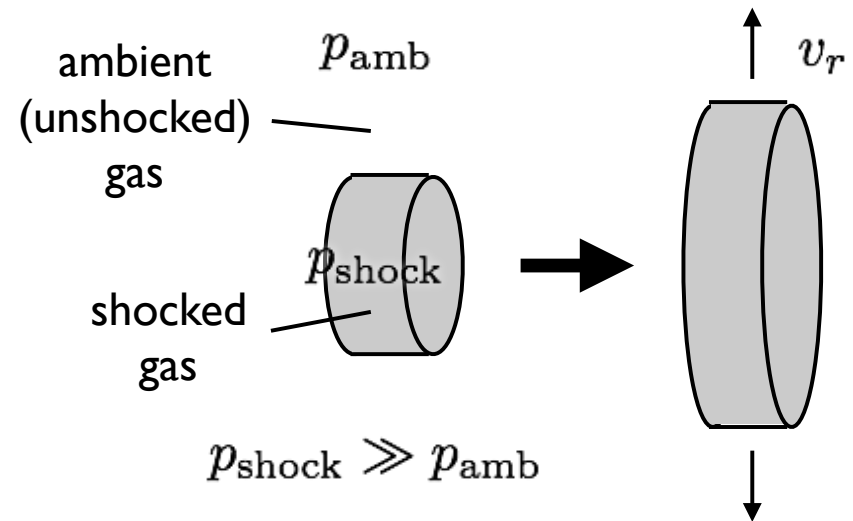
post-shock gas is **strongly** compressed and heated.



oblique shock

Compression and heating are relatively weaker.

adiabatic expansion by pressure gradient
(one-zone model):



timescale of expansion:

$$t_{s0} \sim R_p / c_{s0}$$

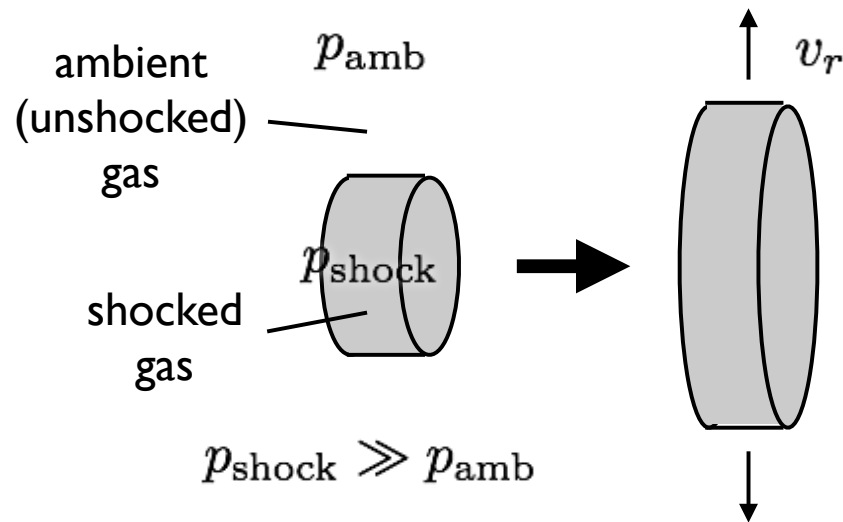
sound speed
($\sim 4 \text{ km s}^{-1}$)

radius of shocked region
 \sim planetesimal radius (assumption)

Shocked gas expands in
sound-crossing time

2. Rapid cooling of silicate vapor

One-zone model



Eq. of motion for vertical direction:

$$\frac{dv_r}{dt} = -\frac{1}{\rho} \frac{\partial p}{\partial r},$$

One-zone approximation:

$$v_r \sim \frac{dR}{dt}, \quad \frac{\partial p}{\partial r} \sim -\frac{p}{R},$$

↓

Eq. of expansion:

$$\frac{d^2 \tilde{R}}{d\tilde{t}^2} = \frac{1}{2} \tilde{R}^{-2\gamma+1},$$

with normalization as

- radius: $\tilde{R} = R/R_p$
- time: $\tilde{t} = t/t_{s0}$
- velocity: $\tilde{v}_r = v_r/c_{c0}$

solution:

- expansion velocity

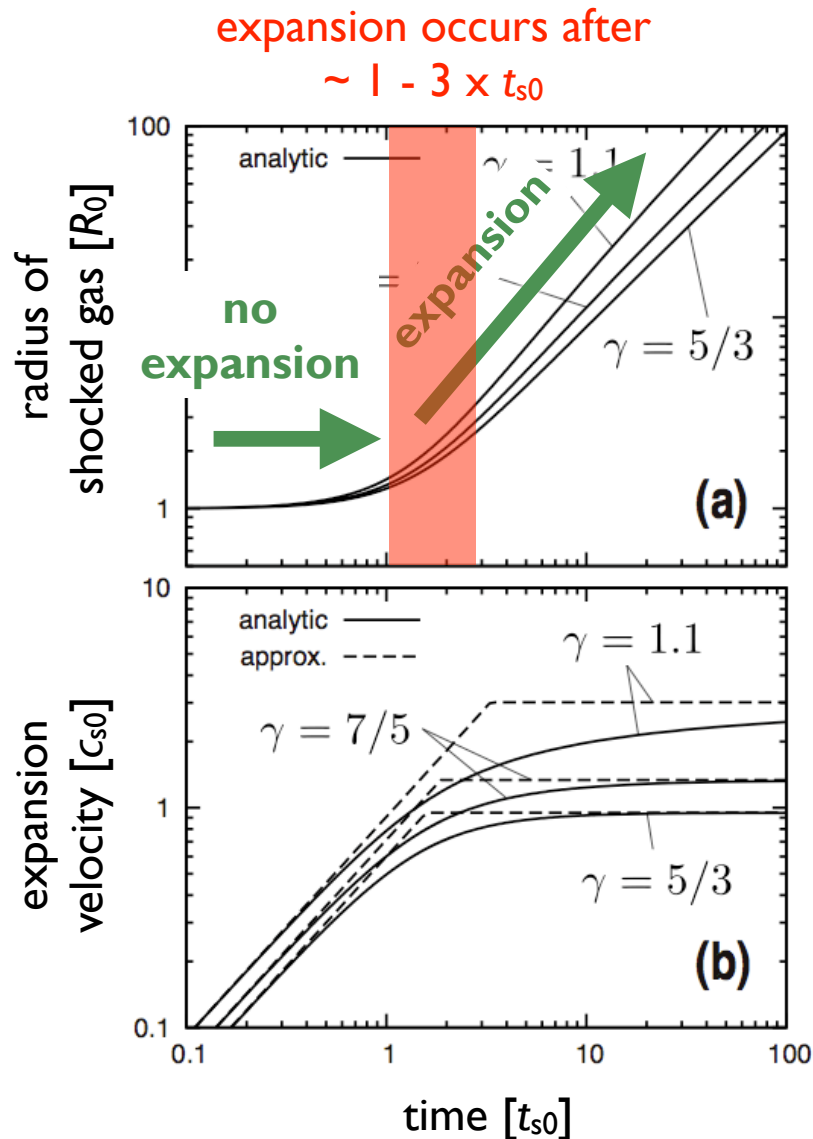
$$\tilde{v}_r = \frac{d\tilde{R}}{d\tilde{t}} = \left[\frac{1 - \tilde{R}^{-2(\gamma-1)}}{\gamma(\gamma-1)} \right]^{1/2}$$

- radius of shocked gas

$$\frac{\tilde{t}}{\sqrt{\gamma(\gamma-1)}} = \int_1^{\tilde{R}} \frac{dy}{\sqrt{1 - y^{-2(\gamma-1)}}}$$

2. Rapid cooling of silicate vapor

Analytic solution



cooling rate of silicate vapor:

$$\begin{aligned}
 - \left(\frac{dT}{dt} \right) &\simeq (0.25 - 0.35) \times T_0 / t_{s0} \\
 &\simeq 2000 \left(\frac{R_p}{1 \text{ km}} \right)^{-1} \left(\frac{T_0}{2000 \text{ K}} \right) \\
 &\quad \times \left(\frac{c_{s0}}{3.7 \text{ km s}^{-1}} \right) \text{ K s}^{-1}
 \end{aligned}$$

- **small** planetesimal \rightarrow **rapid** cooling
- **large** planetesimal \rightarrow **slower** cooling

“rapid cooling of silicate vapor”

3. Condensation of cosmic crystals

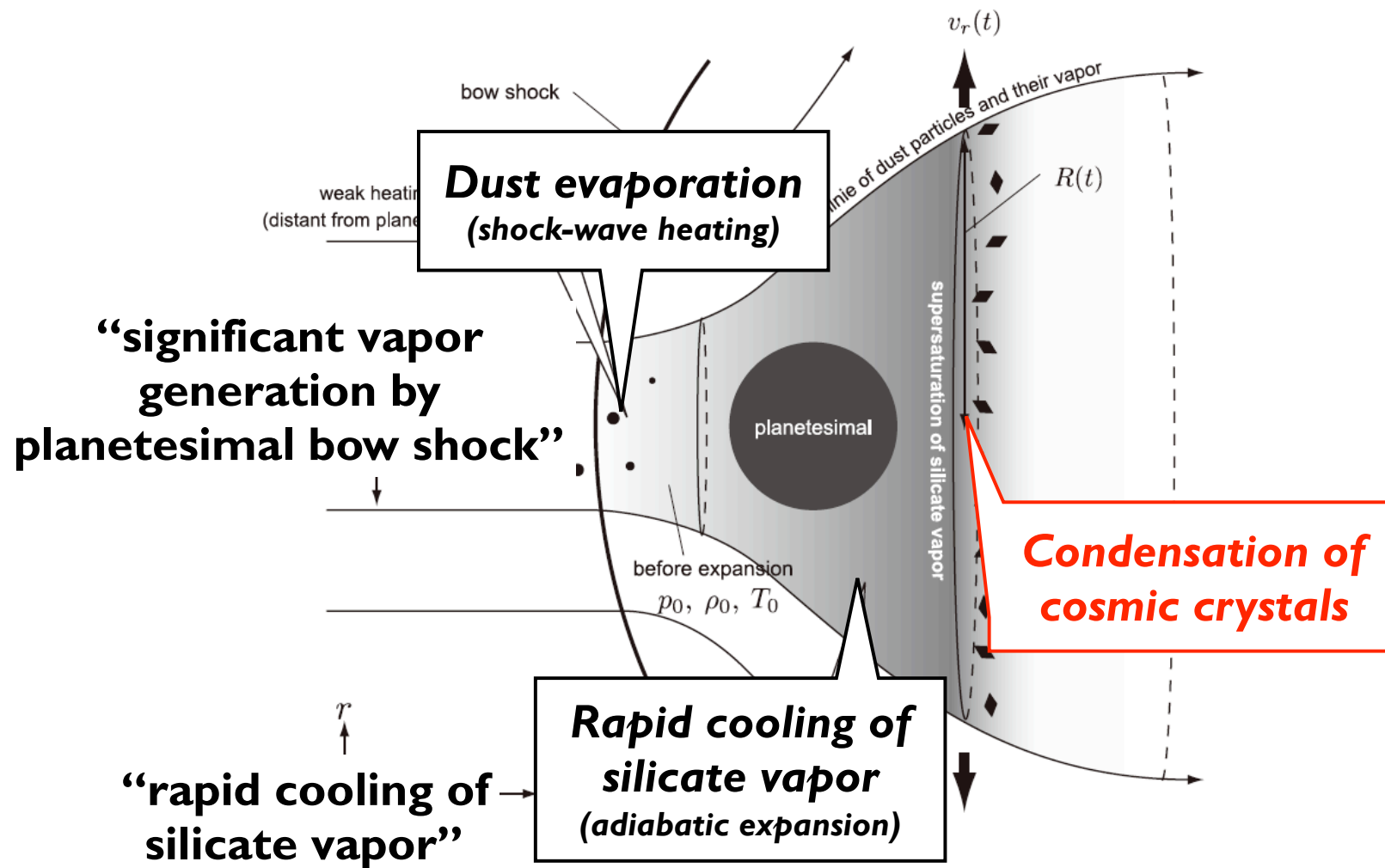
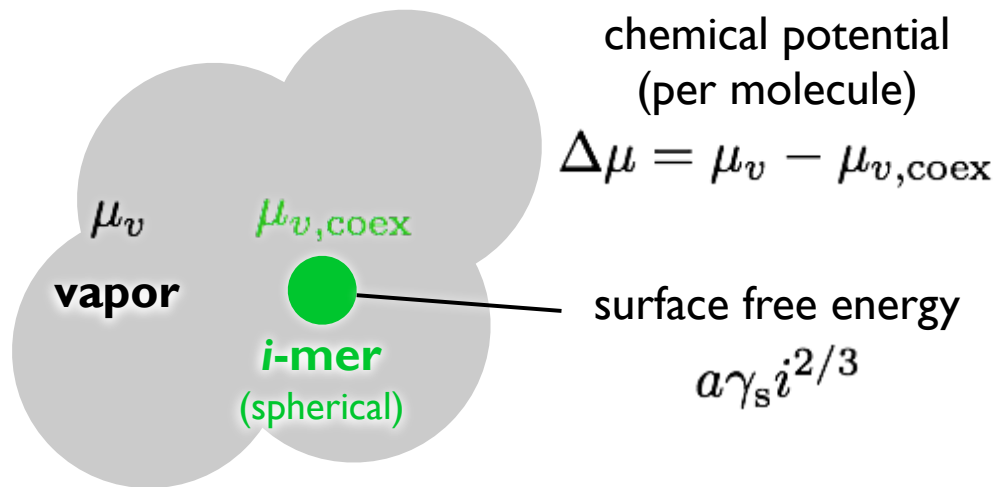


Fig. schematic of planetesimal bow shock
(Miura+2010,ApJ 719, 642)

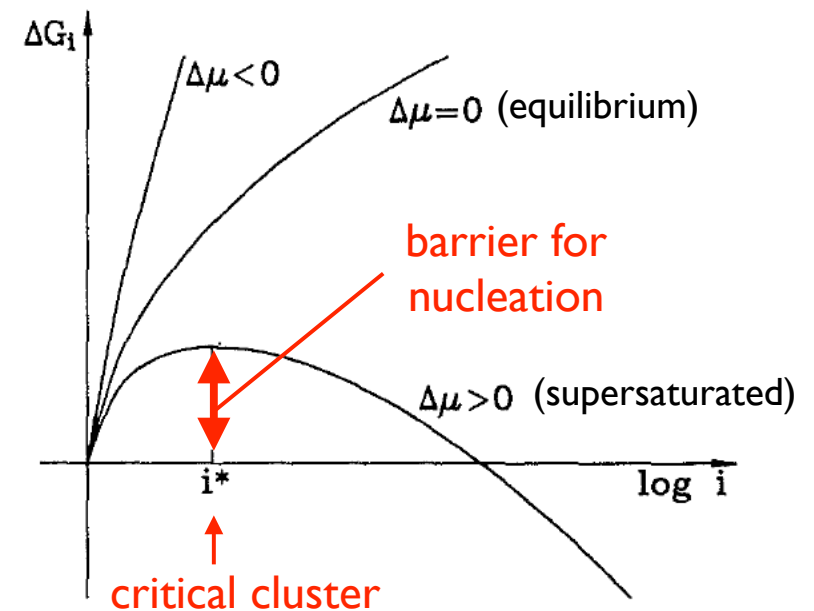
3. Condensation of cosmic crystals

Homogeneous nucleation



Gibbs free energy of formation:

$$\Delta G_i = a\gamma_s i^{2/3} - i\Delta\mu$$



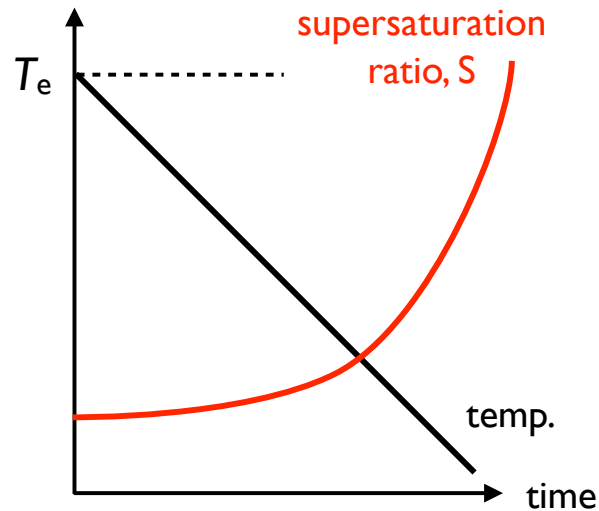
Feder+1996, Adv. Phys. 15, 111
 Dillmann and Meier 1991, J. Chem. Phys. 94, 3872

“delay of nucleation by surface free energy”

3. Condensation of cosmic crystals

Cooling parameter Λ

nucleation and growth in monotonically cooling gas
(Yamamoto and Hasegawa 1977, Prog.Theo. Phys. 58, 816)



$$S(t) = \frac{c_1(t)}{c_1(0)} \exp\left(\frac{t}{\tau_{\text{sat}}}\right)$$

depletion of gas number density timescale of cooling (increase of S)

Only two non-dimensional parameters determine

- (actual) condensation temperature,
- size distribution of condensed grains.

Cooling timescale:

$$\Lambda = \tau_{\text{sat}} / \tau_{\text{coll}}$$

collision interval of vapor molecules

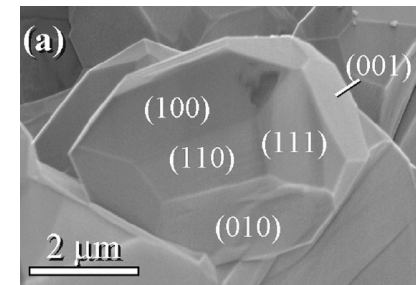
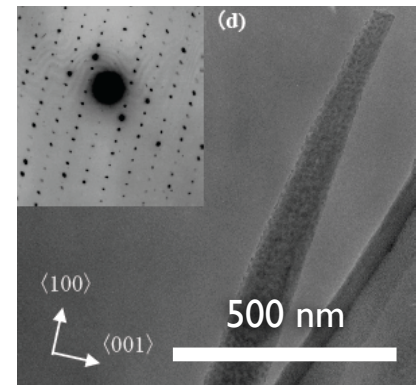
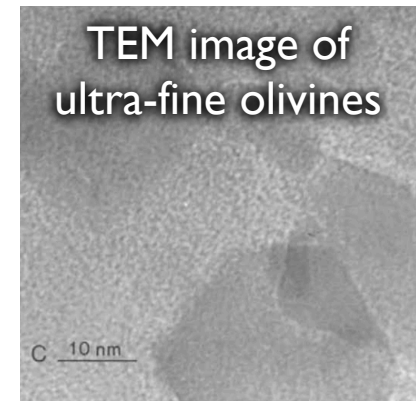
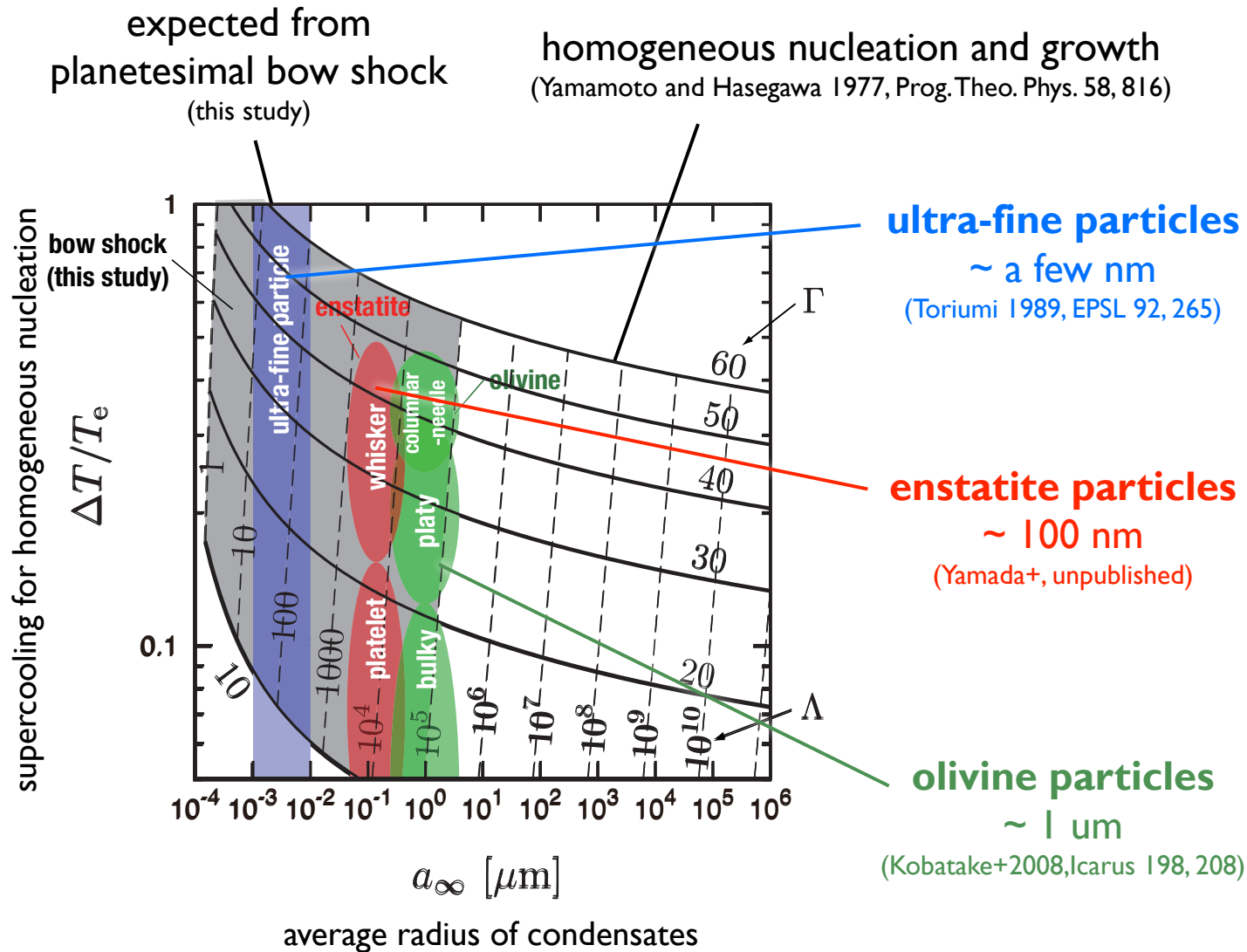
Surface energy of a vapor molecule:

$$\Gamma = \frac{4\pi a_0^2 \gamma_s}{k_B T_e}$$

monomer radius condensation temp. in equilibrium

3. Condensation of cosmic crystals

Diagram of condensed particles



3. Condensation of cosmic crystals

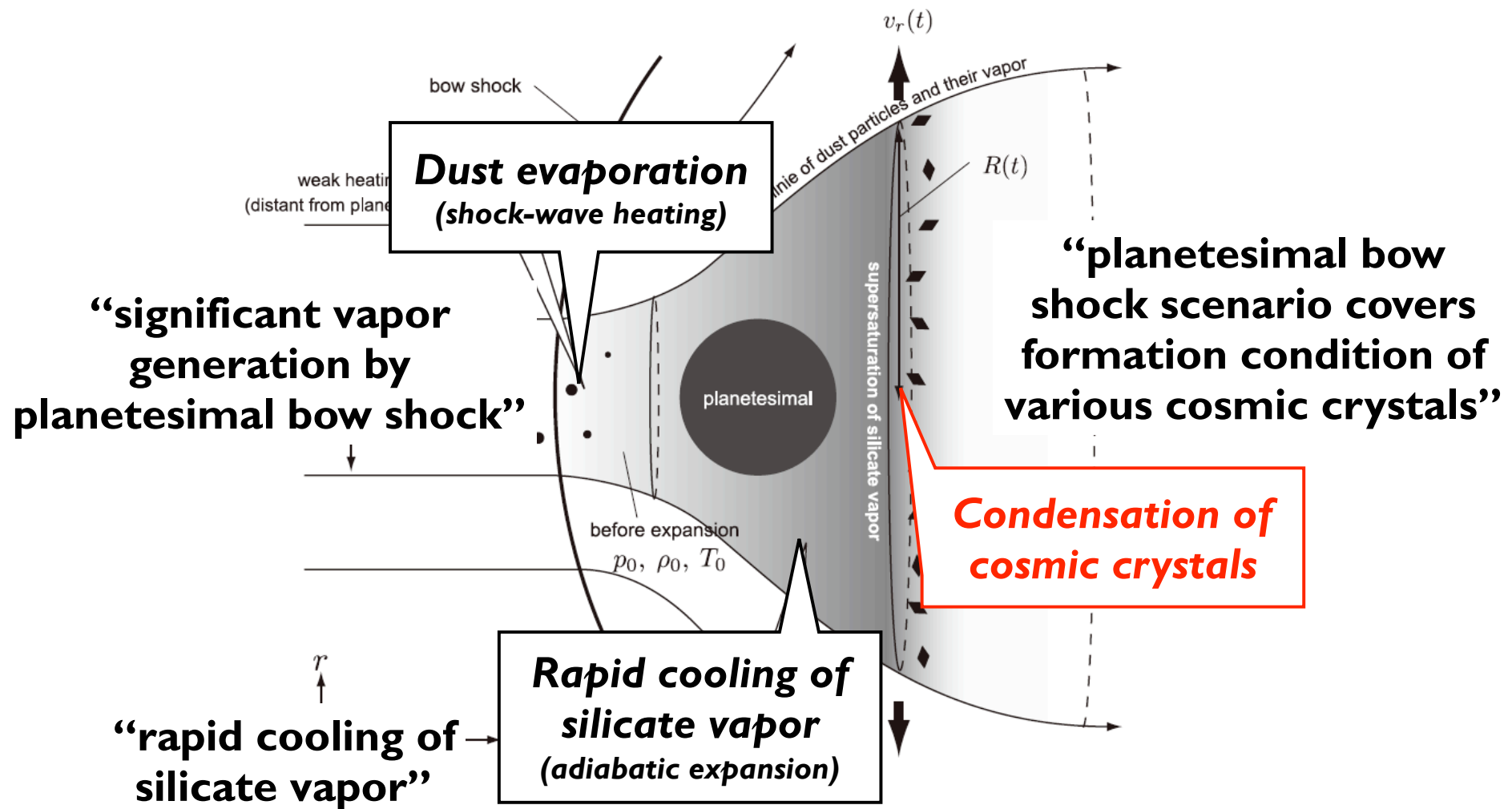


Fig. schematic of planetesimal bow shock
(Miura+2010,ApJ 719, 642)

Conclusions

- Dust evaporation and condensation experiments showed that cosmic crystals with various morphologies were formed from highly-supercooled (supersaturated) silicate vapor. The morphology depends on temperature and supercooling (morphodrom).
- Planetesimal bow shock is one of the candidates for the cosmic crystal formation. It evaporates μm -sized fine silicate particles behind the shock front. The silicate vapor cools rapidly due to the adiabatic expansion.
- Depending on the shock conditions (planetesimal radius, shock velocity, gas number density, and dust-to-gas mass ratio), variety of cosmic crystals in sizes (nm-size to μm -size) and morphology (bulky, platy, whisker, and so forth) was produced.